

Dirac's bra-ket notation

Dirac's bra-ket notation is frequently used to denote vectors in quantum mechanics. As it is not difficult, you will benefit by learning it; it will facilitate reading quantum mechanics books. This article was written for students who have some understandings of linear algebra.

There are two kinds of vectors in bra-ket notation: bra vectors and ket vectors. A bra vector $\langle v|$ denotes a row vector, such as $(v_x \ v_y \ v_z)$. A ket vector $|v\rangle$ denotes a column vector such as $\begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix}$. Therefore a bra vector and a ket vector are dual to each other; (i.e. they can be paired together to produce a scalar) I will show you shortly how one can denote a dot product between vector \vec{A} and vector \vec{B} in bra-ket notation. Let's first denote this product in matrix notation. It would be

$$(A_x \ A_y \ A_z) \begin{pmatrix} B_x \\ B_y \\ B_z \end{pmatrix} = A_x B_x + A_y B_y + A_z B_z \quad (1)$$

. You can easily see that the row vector is on the left side and the column vector is on the right side. Therefore, you have to write a bra vector on the left, and a ket vector on the right. Then it would be $\langle A|B\rangle$. Of course, you may wonder why it has not been written as $\langle A||B\rangle$ but, it's just a convention that you write only one bar in the middle. Now, I will show you how you can write the completeness relation in terms of bra-ket notation. A completeness relation in three dimensions in matrix notation is as follows.

$$\begin{aligned} I &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} (1 \ 0 \ 0) + \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} (0 \ 1 \ 0) + \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} (0 \ 0 \ 1) \end{aligned}$$

You may wonder how this sum could be an identity matrix.

You can easily see it, by multiplying by a row vector $(v_x \ v_y \ v_z)$ on the left-hand side of the sum. Then you get

$$v_x (1 \ 0 \ 0) + v_y (0 \ 1 \ 0) + v_z (0 \ 0 \ 1) = (v_x \ v_y \ v_z) \quad (2)$$

Similarly, you can multiply by a column vector on the right-hand side of sum, and you get the original column vector. Therefore, it is easy to see that this sum is an identity matrix. Here, you can see that the column vectors are on the left side of the row vectors. Therefore, ket vectors precede bra vectors in the completeness relation. So, you can write this as

$$I = \sum_{n=1}^3 |e_n\rangle\langle e_n| \quad (3)$$

in bra-ket notation, where $\begin{pmatrix} 1 & 0 & 0 \end{pmatrix}$ denotes $\langle e_1|$, $\begin{pmatrix} 0 & 1 & 0 \end{pmatrix}$ denotes $\langle e_2|$, and $\begin{pmatrix} 0 & 0 & 1 \end{pmatrix}$ denotes $\langle e_3|$.

Physicists frequently use the completeness relation in quantum mechanics. When there are infinite bases, and the index used to label bases is continuous (in above case, this index was n), the sum above is replaced by an integration. The most common examples are the position matrix x and the momentum matrix p . The completeness relation in these cases can be written as follows.

$$I = \int dx |x\rangle\langle x| = \int dp |p\rangle\langle p| \quad (4)$$

Here, $|x\rangle$ denotes the suitably normalized eigenvector with eigenvalue x . Similarly, $|p\rangle$ denotes the suitably normalized eigenvector with eigenvalue p . Here, I used the word “suitably” to reflect the fact that the eigenvectors with continuous eigenvalues must be normalized differently than the ones with discrete eigenvalues.¹

Problem 1. Using our earlier notations for $|e_n\rangle$ and $\langle e_n|$, convince yourself of the following:

$$\langle e_i|H|e_j\rangle = H_{ij} \quad (5)$$

where H is a matrix and H_{ij} its components.

Summary

- $\langle v|$ is a bra vector, and $|v\rangle$ is a ket vector.
- An inner product between a bra vector $\langle u|$ and a ket vector $|v\rangle$ is denoted as $\langle u|v\rangle$.
- The completeness relation is given by $I = \sum_i |e_i\rangle\langle e_i| = \int dx |x\rangle\langle x| = \int dp |p\rangle\langle p|$.

¹In other words, while the norm of $|e_n\rangle$ can be normalized to be 1 for any n , the norm of $|x\rangle$ can never be so for any x . We will obtain the suitable norm of $|x\rangle$ in our later article “A short introduction to quantum mechanics VI: position basis and the Dirac delta function.”