

Non-Abelian gauge theory in differential forms

In this article, we introduce how non-Abelian gauge theory can be expressed in the language of differential forms often used in mathematics.

Earlier, the covariant derivative was defined as follows:

$$D_\mu \psi = (\partial_\mu - igA_\mu)\psi \quad (1)$$

Here we will re-express the above equation in forms as:

$$D = d + A \quad (2)$$

Here we have suppressed the indices (i.e. μ) and absorbed the factor $-ig$ into the definition of A .

Now $[D_\mu, D_\nu]\psi$ can be expressed as follows:

$$\begin{aligned} D^2\psi &= (d + A)(d + A)\psi = (d + A)(d\psi + A \wedge \psi) \\ &= d^2\psi + A \wedge d\psi + dA \wedge \psi - A \wedge d\psi + A \wedge A\psi \\ &= (dA + A \wedge A)\psi \end{aligned} \quad (3)$$

where from the first line to the second line we have used the fact that A is a one-form and the relation

$$d(\alpha \wedge \beta) = d\alpha \wedge \beta + (-1)^{\deg \alpha}(\alpha \wedge d\beta) \quad (4)$$

Summarizing, we can define the field strength F as follows:

$$F = dA + A \wedge A \quad (5)$$

Also, at this point, it might seem odd that $A \wedge A$ in the above equation is not zero, even though the wedge-product of an object is taken with itself. This term can be expressed in component form as follows:

$$\begin{aligned} A \wedge A &= A_\mu dx^\mu \wedge A_\nu dx^\nu \\ &= A_\mu A_\nu dx^\mu \wedge dx^\nu = A_\nu A_\mu dx^\nu \wedge dx^\mu = -A_\nu A_\mu dx^\mu \wedge dx^\nu \\ &= \frac{A \wedge A}{2} + \frac{A \wedge A}{2} = \frac{A_\mu A_\nu}{2} - \frac{A_\nu A_\mu}{2} dx^\mu dx^\nu \\ &= \frac{1}{2}[A_\mu, A_\nu] dx^\mu \wedge dx^\nu \end{aligned} \quad (6)$$

Using this, we can also express (5) as follows:

$$F = dA + \frac{1}{2}[A, A] \quad (7)$$

We can also derive the equation of motion that F satisfies

$$\begin{aligned}
dF &= d^2 A + dA \wedge A - A \wedge dA = dA \wedge A - A \wedge dA \\
A \wedge F &= A \wedge dA + A \wedge A \wedge A \\
F \wedge A &= dA \wedge A + A \wedge A \wedge A
\end{aligned} \tag{8}$$

Therefore we conclude:

$$\begin{aligned}
dF + A \wedge F - F \wedge A &= 0 \\
dF + [A, F] &= 0
\end{aligned} \tag{9}$$

This is called the Bianchi identity. The last step can be shown as follows:

$$A \wedge F - F \wedge A = A_\lambda \left(\frac{1}{2} F_{\mu\nu} \right) dx^\lambda \wedge dx^\mu \wedge dx^\nu - \left(\frac{1}{2} F_{\mu\nu} \right) A_\lambda dx^\mu \wedge dx^\nu \wedge dx^\lambda \tag{10}$$

$$= [A_\lambda, \frac{1}{2} F_{\mu\nu}] dx^\lambda \wedge dx^\mu \wedge dx^\nu = [A, F] \tag{11}$$

Also, if we obtain the equation of motion by varying the Yang-Mills action (i.e. $\text{Tr} F_{\mu\nu} F^{\mu\nu}$), we obtain:

$$d * F + [A, *F] = 0 \tag{12}$$

We will not show the proof. In the presence of charge, as we considered in our earlier article “The Maxwell Lagrangian and its equation of motion,” the right-hand side of (12) (i.e. 0) is simply replaced by J .

Finally, in the earlier article “Non-Abelian gauge theory,” we mentioned the following infinitesimal gauge transformation:

$$A_\mu^a \rightarrow A_\mu^a + \partial_\mu \theta^a - g f^{abc} \theta^b A_\mu^c \tag{13}$$

$$F_{\mu\nu}^a \rightarrow F_{\mu\nu}^a - g f^{abc} \theta^b F_{\mu\nu}^c \tag{14}$$

In the differential forms language these are translated into:

$$A \rightarrow A + d\theta + [A, \theta] \tag{15}$$

$$F \rightarrow F + [F, \theta] \tag{16}$$

These equations can be shown by using following facts:

$$A = -ig A_\mu dx^\mu = -ig A_\mu^a T^a dx^\mu \tag{17}$$

$$\theta = -ig \theta^b T^b \tag{18}$$

$$F = (-ig) \frac{1}{2} F_{\mu\nu} dx^\mu \wedge dx^\nu = (-ig) \frac{1}{2} F_{\mu\nu}^a T^a dx^\mu \wedge dx^\nu \tag{19}$$

$$[T^a, T^b] = i f^{abc} T^c \tag{20}$$

where we have revived the $-ig$ factor that we had absorbed in the definition of A in (2).

Notice that (15) is exactly $\omega^a \rightarrow \omega^a + D\Lambda^a$ obtained in “Gauge transformation in dreibein.”

Problem 1. Convince yourself that (12) is equivalent to $D_\mu F^{\mu\nu} = 0$ obtained in our earlier article.

Summary

- In differential form notation, a covariant derivative is often written as

$$D = d + A$$

- Then, the field strength is given by

$$D^2 = F = dA + A \wedge A$$

- The Bianchi identity is given by

$$dF + A \wedge F - F \wedge A = dF + [A, F] = 0$$

- The equation of motion is given by

$$d * F + [A, *F] = 0$$

- Under the gauge transformation, we have

$$A \rightarrow A + D\theta$$

$$F \rightarrow F + [F, \theta]$$