

# A short introduction to quantum mechanics X: position and momentum basis and Fourier transformation

We have already noted that  $P$ , the momentum operator, acts by  $-i\hbar\frac{\partial}{\partial x}$ . (Some textbooks use  $\hat{p}$  for the momentum operator.) In other words, if we have:

$$|\psi\rangle = \int dx |x\rangle \langle x|\psi\rangle = \int \psi(x) |x\rangle dx \quad (1)$$

then it follows that:

$$P|\psi\rangle = \int \left( -i\hbar \frac{\partial \psi(x)}{\partial x} \right) |x\rangle dx = \int dx |x\rangle \langle x|P|\psi\rangle \quad (2)$$

Therefore, we have:

$$-i\hbar \frac{\partial}{\partial x} \langle x|\psi\rangle = \langle x|P|\psi\rangle \quad (3)$$

Now, let's plug in  $|p\rangle$ , the eigenvector of the momentum matrix with eigenvalue  $p$ , for  $|\psi\rangle$ . We get:

$$\begin{aligned} -i\hbar \frac{\partial}{\partial x} \langle x|p\rangle &= \langle x|P|p\rangle \\ -i\hbar \frac{\partial}{\partial x} \langle x|p\rangle &= \langle x|p\rangle p \\ -i\hbar \frac{\partial}{\partial x} \langle x|p\rangle &= p \langle x|p\rangle \end{aligned} \quad (4)$$

To transition from the second line to the third line, we used the fact that  $p$  is merely a number, since it's an eigenvalue.

So, this is a differential equation for  $\langle x|p\rangle$  and the solution is given by:

$$\langle x|p\rangle = C e^{ipx/\hbar} \quad (5)$$

Then we can say:

$$|p\rangle = \int dx |x\rangle \langle x|p\rangle = \int dx |x\rangle C e^{ipx/\hbar} \quad (6)$$

Given this, I will now explain the relation between wave functions written in the position basis and those written in the momentum basis. We may write a wave function  $|\psi\rangle$  as follows:

$$|\psi\rangle = \int dx |x\rangle \langle x|\psi\rangle = \int dx \psi(x) |x\rangle \quad (7)$$

In the momentum basis, the same wave function  $|\psi\rangle$  can be expressed as follows:

$$|\psi\rangle = \int dp |p\rangle \langle p|\psi\rangle = \int dp \phi(p) |p\rangle \quad (8)$$

where  $\phi(p) = \langle p|\psi\rangle$

Now, observe:

$$\langle x|\psi\rangle = \int \langle x|p\rangle \langle p|\psi\rangle dp \quad (9)$$

$$\psi(x) = \int C \phi(p) e^{ipx/\hbar} dp \quad (10)$$

Given the above formula, could we express  $\phi(p)$  in terms of  $\psi(x)$ ? If you are careful enough, you will see that it is just a Fourier transformation problem. So, roughly speaking, we should have,  $\phi(p) \sim \int \psi(x) e^{-ipx/\hbar} dx$  with an overall factor to be determined. Remarkably, it is possible to reproduce this as follows:

$$\langle p|\psi\rangle = \int \langle p|x\rangle \langle x|\psi\rangle dx \quad (11)$$

$$\phi(p) = \int C^* \psi(x) e^{-ipx/\hbar} dx \quad (12)$$

where we have used the fact that  $\langle p|x\rangle = \langle x|p\rangle^*$ . So, we have reproduced the so-called inverse Fourier transformation as advertised!

By comparing (10) and (12) with (6) and (7) of my article “Fourier transformations” and assuming that  $C$  is real without a loss of generality, we can get  $C = 1/(\sqrt{2\pi\hbar})$ . The fact that  $\langle x|p\rangle$  is proportional to  $e^{ipx/\hbar}$  makes sense, as it leads to the correct Fourier transformations; we see that quantum mechanics is mathematically consistent. If you study physics and math further, you will frequently encounter such beautiful consistencies!

To summarize, we derived:

$$\psi(x) = \int \frac{dp}{\sqrt{2\pi\hbar}} \phi(p) e^{ipx/\hbar} \quad (13)$$

$$\phi(p) = \int \frac{dx}{\sqrt{2\pi\hbar}} \psi(x) e^{-ipx/\hbar} \quad (14)$$

So far, in position basis, we have seen that  $X$  acts by multiplying the position-space wave function by  $x$  and  $P$  acts by  $-i\hbar \frac{\partial}{\partial x}$  where the position-space wave function for a state  $|\psi\rangle$  is given by  $\langle x|\psi\rangle$ . Then, a natural

question to ask is how  $X$  and  $P$  act on the momentum-space wave function  $\langle p|\psi\rangle$ .

How  $P$  acts is easy. We have:

$$\langle p|P|\psi\rangle = (\langle p|P)|\psi\rangle = p\langle p|\psi\rangle \quad (15)$$

where we have used  $P|p\rangle = p|p\rangle$ . In other words,  $P$  acts by multiplying by  $p$ .

To figure out how  $X$  acts, let's consider the following equation. (**Problem 1.** Prove this! Hint: Insert  $I = \int dx|x\rangle\langle x|$  between  $e^{iaX}$  and  $|p\rangle$ )

$$e^{iaX}|p\rangle = |p + \hbar a\rangle \quad (16)$$

For infinitesimal  $a = \Delta p/\hbar$ , the above equation implies

$$\begin{aligned} e^{i\Delta p X/\hbar}|p\rangle &= |p + \Delta p\rangle \\ \left(1 + \frac{i\Delta p X}{\hbar}\right)|p\rangle &= |p + \Delta p\rangle \\ \frac{i\Delta p X}{\hbar}|p\rangle &= |p + \Delta p\rangle - |p\rangle \end{aligned} \quad (17)$$

Given this, we have:

$$\begin{aligned} |\psi\rangle &= \int dp|p\rangle\langle p|\psi\rangle = \int dp|p\rangle\psi(p) \\ \frac{i\Delta p X}{\hbar}|\psi\rangle &= \int dp \frac{i\Delta p X}{\hbar}|p\rangle\langle p|\psi\rangle \\ &= \int dp(|p + \Delta p\rangle - |p\rangle)\langle p|\psi\rangle \\ &= \int dp|p + \Delta p\rangle\langle p|\psi\rangle - \int dp|p\rangle\langle p|\psi\rangle \\ &= \int dp|p\rangle\langle p - \Delta p|\psi\rangle - \int dp|p\rangle\langle p|\psi\rangle \\ &= \int dp(\psi(p - \Delta p) - \psi(p))|p\rangle \\ \frac{i\Delta p X}{\hbar}|\psi\rangle &= \int dp \left(-\frac{\partial\psi}{\partial p}\Delta p\right)|p\rangle \\ X|\psi\rangle &= \int dp \left(i\hbar\frac{\partial\psi}{\partial p}\right)|p\rangle \end{aligned} \quad (18)$$

Therefore, we see that, in the momentum basis, the position operator  $X$  acts by  $i\hbar\frac{\partial}{\partial p}$ . Summarizing, we have:

$$X\psi(p) = i\hbar\frac{\partial\psi(p)}{\partial p} \quad (19)$$

$$P\psi(p) = p\psi(p) \quad (20)$$

where  $\psi(p) = \langle p|\psi\rangle$ .

**Problem 2.** Check  $XP - PX = i\hbar$  in momentum basis.

Compare (19) and (20) with the following formulas for the familiar position basis:

$$X\psi(x) = x\psi(x) \quad (21)$$

$$P\psi(x) = -i\hbar \frac{\partial\psi(x)}{\partial x} \quad (22)$$

where  $\psi(x) = \langle x|\psi\rangle$ . We see that there is a symmetry upon simultaneous exchange of  $x$  for  $p$  and of  $i$  for  $-i$ .

Finally, let me conclude this article with some observations on Dirac delta function. Changing the integration variable from  $x$  to  $x'$  in (14), and plugging this to (13), we obtain:

$$\psi(x) = \int dx' \psi(x') \left( \int \frac{dp}{2\pi\hbar} e^{ip(x-x')/\hbar} \right) \quad (23)$$

Therefore, we conclude:

$$\delta(x - x') = \int \frac{dp}{2\pi\hbar} e^{ip(x-x')/\hbar} \quad (24)$$

Therefore, we see that the condition that one must come back to the original function, if one Fourier-transform and inverse-Fourier-transform it, yields a formula for Dirac delta function.

In other words, we derived an explicit formula for delta function as follows:

$$\delta(x) = \int \frac{dp}{2\pi} e^{ipx} \quad (25)$$

Actually, we can derive this equation more easily as follows:

$$\begin{aligned} \delta(x - x') &= \langle x|x'\rangle = \int dp \langle x|p\rangle \langle p|x'\rangle \\ &= \int \frac{dp}{2\pi\hbar} e^{ip(x-x')/\hbar} \end{aligned} \quad (26)$$

Now more problems. Consider a wave function  $\phi(x)$  as follows:

$$\phi(x) = C_1 \exp \left[ ikx - \frac{x^2}{2d^2} \right] \quad (27)$$

**Problem 3.** Use Taylor series to show

$$e^{iPa/\hbar} \psi(x) = \psi(x + a) \quad (28)$$

In other words, the momentum operator  $P$  moves the position of the wave function by a certain amount. Thus, we say  $P$  generates “translation” (i.e.

the movement of position). In the next article, we will see that the Hamiltonian operator  $H$  moves the time of the wave function. In other words, we say  $H$  generates “time evolution.”

**Problem 4.** Show that the probability density function for position due to (27) is given by a normal (i.e. a Gaussian) distribution. (Hint: See our earlier article “Probability density function.”) Thus, we call such a wave “Gaussian wave packet.”

**Problem 5.** Obtain,  $C_1$  by assuming that  $\phi(x)$  is properly normalized.

**Problem 6.** Find  $\Delta x$ , the standard deviation of the position  $x$ , in terms of  $d$ . Find the expectation value of the momentum  $p$  without Fourier-transforming  $\phi(x)$ . Find the expectation value of  $p^2$  similarly. Thus, obtain  $\Delta p$ , the standard deviation of the momentum.

**Problem 7.** Fourier transform  $\phi(x)$  to obtain  $\psi(p)$ , the wave function in momentum space. Thus, show that the probability density function in the momentum space is also given by a normal distribution. From this, obtain the expectation value of  $p$  and the standard deviation  $\Delta p$  again. Check that they agree with the values you obtained using the wave function in position basis. Check also that following holds.

$$\Delta x \Delta p = \frac{\hbar}{2} \quad (29)$$

This, you found in case of Gaussian packet. Generally, the following holds:

$$\Delta x \Delta p \geq \frac{\hbar}{2} \quad (30)$$

This is called “Heisenberg’s uncertainty principle.” We will talk more about this in a later article.

## Summary

- $\langle x|p\rangle$  is proportional to  $e^{ipx/\hbar}$ .
- The relation between position basis and momentum basis is that of Fourier transformation and inverse Fourier transformation.